

Supplement for Safe Task Space Synchronization with Time-Delayed Information

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In this supplement, bounds on certain terms of the Lyapunov function derivative are computed. The results of these bounds are used in the main paper.

APPENDIX A DEVELOPMENT OF A BOUNDS ON \dot{P}_1 , \dot{P}_2 AND \dot{P}_3

The time derivatives of the LK functional terms using Leibniz rule yield

$$\dot{P}_1 \leq T c_1 - \frac{K_{LK_1} \omega_1}{2} \int_{t_T}^t \ddot{p}_h^T(l) \ddot{p}_h(l) dl$$

where $c_1 \geq \frac{K_{LK_1} \omega_1}{2} \|\ddot{p}_h(t)\|^2 \in \mathbb{R}^+$ is a constant. Similarly, the derivative of P_2 can be computed as follows

$$\begin{aligned} \dot{P}_2 &= \frac{K_{LK_2} \omega_2 T}{2} \|\dot{\phi}(e_p) e_p(t) + \phi(e_p)(t) \dot{e}_p(t)\|^2 \\ &\quad - \frac{K_{LK_2} \omega_2}{2} \int_{t_T}^t \left[\frac{d}{dt} (\phi(e_p)(l) e_p(l)) \right]^T \left[\frac{d}{dt} (\phi(e_p)(l) e_p(l)) \right] dl \end{aligned}$$

where $\phi(e_p)$, and $\dot{\phi}(e_p)$ denote diagonal matrices containing terms $\phi(e_{pi})$, and $\dot{\phi}(e_{pi})$, $\forall i = \{1, \dots, n\}$, respectively. Substituting for $\dot{\phi}$, $\dot{\eta}$ and \dot{e}_p , using triangle and Young's inequalities, the following upper bound on the first term of \dot{P}_2 is obtained

$$\begin{aligned} &k_1^2 6 K_{LK_2} \omega_2 T \|\phi^3(e_p) e_p^2\|^2 \|e_p\|^2 \\ &+ 6 K_{LK_2} \omega_2 T w_j^2 \|\phi^2(e_p) e_p^2\|^2 \|\dot{\theta}\|^2 \|\tilde{\zeta}_j\|^2 \\ &+ 6 K_{LK_2} \omega_2 T j_p^2 \|\phi^2(e_p) e_p^2\|^2 \|\eta\|^2 \\ &+ k_1^2 6 K_{LK_2} \omega_2 T \|\phi^2(e_p) e_p\|^2 \\ &+ 6 K_{LK_2} \omega_2 T w_j^2 \|\phi(e_p)\|^2 \|\dot{\theta}\|^2 \|\tilde{\zeta}_j\|^2 \\ &+ 6 K_{LK_2} \omega_2 T j_p^2 \|\phi(e_p)\|^2 \|\eta\|^2 \leq T \rho_1(\|y\|) \|y\|^2 \quad (1) \end{aligned}$$

where $y = [e_p^T \ \eta^T \ \tilde{\zeta}_j^T]^T$, $\rho_1(\|y\|)$ is a positive invertible non-decreasing function. The symbol $\phi^3(e_p) e_p^2 = [\phi^3(e_{p1}) e_{p1}^2, \dots, \phi^3(e_{pn}) e_{pn}^2]^T$ is a vector containing elements $\phi^3(e_{pi}) e_{pi}^2$, $\phi^2(e_{pi}) e_{pi}^2 = [\phi^2(e_{p1}) e_{p1}^2, \dots, \phi^2(e_{pn}) e_{pn}^2]^T$ is a vector containing elements of $\phi^2(e_{pi}) e_{pi}^2$, and $\phi^2(e_p) e_p$ is a vector containing elements $\phi^2(e_{pi}) e_{pi}$.

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\dot{P}_3 is computed as

$$\dot{P}_3 \leq T c_2 - \frac{K_{LK_3} \omega_3}{2} \int_{t_T}^t \dot{p}_h^T(l) \dot{p}_h(l) dl$$

where $c_2 \geq \frac{K_{LK_3} \omega_3}{2} \|\dot{p}_h(t)\|^2 \in \mathbb{R}^+$ is a constant.

APPENDIX B

DEVELOPMENT OF A BOUND ON $\eta^T (W_y \zeta - W_{yT} \zeta)$

Consider following term from (34) (of the main paper)

$$\begin{aligned} \eta^T (W_y - W_{yT}) \zeta_y &= M[\dot{j}_p^{-1}(\dot{p}_h - \dot{p}_{hT}) \\ &+ k_1 \dot{j}_p^{-1}(\phi(e_p) e_p - \phi(e_{pT}) e_{pT}) \\ &+ \hat{j}_p^{-1}(\ddot{p}_h - \ddot{p}_{hT}) + k_1 \hat{j}_p^{-1}(\dot{\phi}(e_p) e_p \\ &- \dot{\phi}(e_{pT}) e_{pT}) + k_1 \hat{j}_p^{-1}(\phi(e_p) \dot{e}_p - \phi(e_{pT}) \dot{e}_{pT})] \\ &+ C(\theta, \dot{\theta})(\eta - \eta_T) \quad (2) \end{aligned}$$

The norm bound can be written as

$$\begin{aligned} \|\eta\| \|(W_y - W_{yT}) \zeta_y\| &\leq \bar{m} \bar{j}_{p_{inv}}^2 \bar{j}_{p_{der}} \|\dot{\theta}\| \|\eta\| \left\| \int_{t_T}^t \ddot{p}_h(l) dl \right\| \\ &+ k_1 \bar{m} \bar{j}_{p_{inv}}^2 \bar{j}_{p_{der}} \|\dot{\theta}\| \|\eta\| \left\| \int_{t_T}^t \frac{d}{dl} \phi(e_p)(l) e_p(l) dl \right\| \\ &+ k_1 \bar{m} \bar{j}_{p_{inv}} \|\eta\| \|\dot{\phi}(e_p) e_p - \dot{\phi}(e_{pT}) e_{pT} + \phi(e_p) \dot{e}_p \\ &- \phi(e_{pT}) \dot{e}_{pT}\| + \bar{c} \|\dot{\theta}\| \|\eta\| \|\eta - \eta_T\| + \bar{m} \bar{j}_{p_{inv}} \bar{p}_3 \|\eta\| \quad (3) \end{aligned}$$

where $\|C(\theta, \dot{\theta})\| \leq \bar{c} \|\dot{\theta}\|$ is used. Using Young's inequality on select terms, and the following bound on $\|\eta - \eta_T\|$ developed using (7) and (9) (of the main paper)

$$\|\eta - \eta_T\|^2 \leq 2 j_{p_i}^2 \left\| \int_{t_T}^t \ddot{p}_h(l) dl \right\|^2 + 2 k_1^2 j_{p_i}^2 \left\| \int_{t_T}^t \frac{d}{dl} (\phi(e_p) e_p) dl \right\|^2$$

the following upper bound can be developed

$$\begin{aligned} \|\eta\| \|(W_y - W_{yT}) \zeta_y\| &\leq \bar{m} \bar{j}_{p_{inv}} \bar{p}_3 \|\eta\| \\ &+ \left(\frac{(k_1 + 1) \bar{m} \bar{j}_{p_{inv}}^2 \bar{j}_{p_{der}} \gamma_7^2 + \bar{c} \gamma_9^2}{4} \|\dot{\theta}\|^2 \right) \|\eta\|^2 \\ &+ \left(\frac{\bar{m} \bar{j}_{p_{inv}}^2 \bar{j}_{p_{der}}}{\gamma_7^2} + \frac{\bar{c} 2 j_{p_i}^2}{\gamma_9^2} \right) \left\| \int_{t_T}^t \ddot{p}_h(l) dl \right\|^2 \\ &+ k_1 \bar{m} \bar{j}_{p_{inv}} \|\eta\| \|\dot{\phi}(e_p) e_p - \dot{\phi}(e_{pT}) e_{pT}\| \\ &+ \|\phi(e_p) \dot{e}_p - \phi(e_{pT}) \dot{e}_{pT}\| \\ &+ \left(\frac{k_1 \bar{m} \bar{j}_{p_{inv}}^2 \bar{j}_{p_{der}}}{\gamma_7^2} + \frac{\bar{c} 2 k_1^2 j_{p_i}^2}{\gamma_9^2} \right) \left\| \int_{t_T}^t \frac{d}{dl} (\phi(e_p) e_p) dl \right\|^2 \quad (4) \end{aligned}$$

To further develop the bound in (4), following bounds are developed. Using the triangle inequality of vector norms results in

$$\begin{aligned} \|\dot{\phi}(e_p)e_p - \dot{\phi}(e_{p_T})e_{p_T}\| &\leq \|2k_1(\phi^3(e_p)e_p^2 - \phi^3(e_{p_T})e_{p_T}^2)\| \\ &+ \|2(\phi^2(e_p)e_p - \phi^2(e_{p_T})e_{p_T})\| \|W_j \tilde{\zeta}_j\| \\ &+ \|2\phi^2(e_p)e_p \hat{J}_p \eta - 2\phi^2(e_{p_T})e_{p_T} \hat{J}_p \eta_T\| \end{aligned}$$

where the symbol $\phi^3(e_p)e_p^2 = [\phi^3(e_{p1})e_{p1}^2, \dots, \phi^3(e_{pn})e_{pn}^2]^T$ is a vector containing elements $\phi^3(e_{pi})e_{pi}^2$ and $\phi^2(e_p)e_p$ is a diagonal matrix containing elements $\phi^2(e_{pi})e_{pi}$. Using Appendix A of [1], bounds on each term can be developed as

$$\begin{aligned} \|2k_1(\phi^3(e_p)e_p^2 - \phi^3(e_{p_T})e_{p_T}^2)\| &\leq \rho_2(\|e_p\|) \|p_h - p_{h_T}\| \\ \|2(\phi^2(e_p)e_p - \phi^2(e_{p_T})e_{p_T})\| &\leq \rho_3(\|e_p\|) \|p_h - p_{h_T}\| \\ \|2\phi^2(e_p)e_p \hat{J}_p \eta - 2\phi^2(e_{p_T})e_{p_T} \hat{J}_p \eta_T\| \\ &\leq \rho_4(\|z_1\|) \|z_1(t) - z_1(t-T)\| \end{aligned} \quad (5)$$

where $\rho_i, \forall i = \{2, 3, 4\}$ are positive invertible functions, $z_1(t) = [e_p^T \ \eta^T]^T \in \mathcal{D} \subset \mathbb{R}^{2n}$ and $z_1(t) - z_1(t-T) = [(p_h - p_{h_T})^T \ (\eta - \eta_T)^T]^T$.

$$\begin{aligned} \|\eta\| \|\dot{\phi}(e_p)e_p - \dot{\phi}(e_{p_T})e_{p_T}\| &\leq \|\eta\| \rho_2(\|e_p\|) \|p_h - p_{h_T}\| \\ &+ \|\eta\| \rho_3(\|e_p\|) \|p_h - p_{h_T}\| \|W_j \tilde{\zeta}_j\| + \rho_4(\|z_1\|) \|\eta\| \|\eta - \eta_T\| \\ &+ \rho_4(\|z_1\|) \|\eta\| \|p_h - p_{h_T}\| \end{aligned} \quad (6)$$

Using Young's inequality

$$\begin{aligned} \|\eta\| \|\dot{\phi}(e_p)e_p - \dot{\phi}(e_{p_T})e_{p_T}\| &\leq \rho_2(\|e_p\|) \|\eta\| \|p_h - p_{h_T}\| \\ &+ \rho_4(\|z_1\|) \|\eta\| \|p_h - p_{h_T}\| \\ &+ \frac{\bar{p}_1 w_j \rho_3^2(\|e_p\|) \|\dot{\theta}\|^2 \gamma_{11}^2 + \rho_4^2(\|z_1\|) \gamma_{12}^2}{4} \|\eta\|^2 + \frac{\bar{p}_1 w_j}{\gamma_{11}^2} \|\tilde{\zeta}_j\|^2 \\ &+ \frac{2j_{p_i}^2}{\gamma_{12}^2} \|\int_{t_T}^t \ddot{p}_h(l) dl\|^2 + \frac{2k_1^2 j_{p_i}^2}{\gamma_{12}^2} \|\int_{t_T}^t \frac{d}{dl}(\phi(e_p)(l)e_p(l)) dl\|^2 \end{aligned} \quad (7)$$

Consider following term from (4)

$$\begin{aligned} \|\phi(e_p)\dot{e}_p - \phi(e_{p_T})\dot{e}_{p_T}\| &\leq k_1(\phi^2(e_p)e_p - \phi^2(e_{p_T})e_{p_T}) \\ &+ \|\phi(e_p) - \phi(e_{p_T})\| \|W_j \tilde{\zeta}_j\| + \|\phi(e_p) \hat{J}_p \eta - \phi(e_{p_T}) \hat{J}_p \eta_T\| \end{aligned}$$

where $\phi(e_p)$ is a diagonal matrix containing terms $\phi(e_{pi})$, $\phi^2(e_p)e_p$ is a vector containing terms $\phi^2(e_{pi})e_{pi}$. Using Appendix A of [1], each term can be norm bounded as

$$\begin{aligned} \|k_1(\phi^2(e_p)e_p - \phi^2(e_{p_T})e_{p_T})\| &\leq \rho_5(\|e_p\|) \|p_h - p_{h_T}\| \\ \|(\phi(e_p) - \phi(e_{p_T}))\| &\leq \rho_6(\|e_p\|) \|p_h - p_{h_T}\| \\ \|\phi(e_p) \hat{J}_p \eta - \phi(e_{p_T}) \hat{J}_p \eta_T\| &\leq \rho_7(\|z_1\|) \|z_1(t) - z_1(t-T)\| \end{aligned}$$

such that $\rho_i, \forall i = \{5, 6, 7\}$ are positive invertible functions. The norm bound can be written as

$$\begin{aligned} \|\eta\| \|\phi(e_p)\dot{e}_p - \phi(e_{p_T})\dot{e}_{p_T}\| &\leq \rho_5(\|e_p\|) \|\eta\| \|p_h - p_{h_T}\| \\ &+ \bar{p}_1 w_j \rho_6(\|e_p\|) \|\dot{\theta}\| \|\eta\| \|\tilde{\zeta}_j\| + \rho_7(\|z_1\|) \|\eta\| \|\eta - \eta_T\| \\ &+ \rho_7(\|z_1\|) \|\eta\| \|p_h - p_{h_T}\| \end{aligned} \quad (8)$$

Using Young's Inequality,

$$\begin{aligned} \|\eta\| \|\phi(e_p)\dot{e}_p - \phi(e_{p_T})\dot{e}_{p_T}\| &\leq \rho_5(\|e_p\|) \|\eta\| \|p_h - p_{h_T}\| \\ &+ \rho_7(\|z_1\|) \|\eta\| \|p_h - p_{h_T}\| \\ &+ \frac{\bar{p}_1 w_j \rho_6^2(\|e_p\|) \|\dot{\theta}\|^2 \gamma_{13}^2 + \rho_7^2(\|z_1\|) \gamma_{14}^2}{4} \|\eta\|^2 + \frac{\bar{p}_1 w_j}{\gamma_{13}^2} \|\tilde{\zeta}_j\|^2 \\ &+ \frac{2j_{p_i}^2}{\gamma_{14}^2} \|\int_{t_T}^t \ddot{p}_h(l) dl\|^2 + \frac{2k_1^2 j_{p_i}^2}{\gamma_{14}^2} \|\int_{t_T}^t \frac{d}{dl}(\phi(e_p)(l)e_p(l)) dl\|^2 \end{aligned} \quad (9)$$

Using (7) and (9), the bound in (4) can be written as

$$\begin{aligned} \|\eta\| \|(W_y - W_{y_T})\tilde{\zeta}_j\| &\leq \bar{m} \bar{j}_{p_{inv}} \bar{p}_3 \|\eta\| + \rho_9(\|z_1\|, \|\dot{\theta}\|) \|\eta\|^2 \\ &+ c_{LK_1} \|\int_{t_T}^t \ddot{p}_h(l) dl\|^2 + c_{LK_2} \|\int_{t_T}^t \frac{d}{dl}(\phi(e_p)(l)e_p(l)) dl\|^2 \\ &+ c_{LK_3} \|\int_{t_T}^t \dot{p}_h dl\|^2 + (\frac{\bar{p}_1 w_j}{\gamma_{11}^2} + \frac{\bar{p}_1 w_j}{\gamma_{13}^2}) \|\tilde{\zeta}_j\|^2 \end{aligned} \quad (10)$$

$$\begin{aligned} \text{where } \rho_8(\|z_1\|) &= \rho_2(\|e_p\|) + \rho_5(\|e_p\|) + \\ &+ \rho_4(\|z_1\|) + \rho_7(\|z_1\|) \rho_9(\|z_1\|, \|\dot{\theta}\|) = \\ &(\frac{(k_1+1) \bar{m} \bar{j}_{p_{inv}}^2 \bar{j}_{p_{der}} \gamma_7^2 + \bar{c} \gamma_9^2 + \bar{p}_1 w_j \rho_3^2(\|e_p\|) \|\dot{\theta}\|^2 \gamma_{11}^2 + \rho_4^2(\|z_1\|) \gamma_{12}^2}{\bar{p}_1 w_j \rho_6^2(\|e_p\|) \|\dot{\theta}\|^2 \gamma_{13}^2 + \rho_7^2(\|z_1\|) \gamma_{14}^2}) \|\dot{\theta}\|^2 + \\ &(\frac{\bar{m} \bar{j}_{p_{inv}}^2 \bar{j}_{p_{der}}}{4} + \frac{\bar{c} 2j_{p_i}^2}{\gamma_9^2} + \frac{2j_{p_i}^2}{\gamma_{12}^2} + \frac{2j_{p_i}^2}{\gamma_{14}^2}), c_{LK_1} = \\ &(\frac{\bar{m} \bar{j}_{p_{inv}}^2 \bar{j}_{p_{der}}}{\gamma_7^2} + \frac{\bar{c} 2k_1^2 j_{p_i}^2}{\gamma_9^2} + \frac{2k_1^2 j_{p_i}^2}{\gamma_{12}^2} + \frac{2k_1^2 j_{p_i}^2}{\gamma_{14}^2}), c_{LK_2} = \\ &(\frac{k_1 \bar{m} \bar{j}_{p_{inv}}^2 \bar{j}_{p_{der}}}{\gamma_7^2} + \frac{\bar{c} 2k_1^2 j_{p_i}^2}{\gamma_9^2} + \frac{2k_1^2 j_{p_i}^2}{\gamma_{12}^2} + \frac{2k_1^2 j_{p_i}^2}{\gamma_{14}^2}), c_{LK_3} = \frac{1}{\gamma_{15}^2}. \end{aligned}$$

APPENDIX C DEVELOPMENT OF BOUNDS ON CROSS TERMS

Using Cauchy-Schwarz, and Young's inequality and $\|W_j\| \leq w_j \|\dot{\theta}\|$

$$\begin{aligned} \|\tilde{\zeta}_j^T W_j^T \int_{t_T}^t \frac{d}{dt} \phi(e_p) e_p dl\| &\leq \frac{\gamma_6^2}{4} w_j^2 \|\dot{\theta}\|^2 \|\tilde{\zeta}_j\|^2 \\ &+ \frac{1}{\gamma_6^2} \|\int_{t_T}^t \frac{d}{dt} \phi(e_p) e_p dl\|^2 \end{aligned} \quad (11)$$

where $\gamma_6 \in \mathbb{R}^+$ is a constant. Similarly using Cauchy-Schwarz, Young's inequality and $\|W_{y_T}\| \leq \rho_2(\|y_T\|) = w_{y_T} \|y_T\| + \bar{w}_{y_T}$, where $y_T = [e_{p_T}^T \ \eta_T^T \ \tilde{\zeta}_j^T]^T \in \mathcal{D} \subset \mathbb{R}^{2n+m_1}$ and $\|W_y\| \leq w_y \|y\| + \bar{w}_y$, we have

$$\begin{aligned} \|\tilde{\zeta}_y^T W_{y_T}^T (\eta - \eta_T)\| &\leq \frac{\gamma_5^2}{4} \|\tilde{\zeta}_y\|^2 + \frac{2j_{p_i}^2}{\gamma_5^2} \rho_{10}^2(\|\eta_T\|) \\ &+ \|\int_{t_T}^t \ddot{p}_h(l) dl\|^2 + \frac{2k_1^2 j_{p_i}^2}{\gamma_5^2} \rho_{10}^2(\|\eta_T\|) \\ &+ \|\int_{t_T}^t \frac{d}{dt}(\phi(e_p)(l)e_p(l)) dl\|^2 \end{aligned} \quad (12)$$

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